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## **Abstract Book**

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### THE EFFECT OF PLASTIC ANISOTROPY AND TENSION COMPRESSION–ASYMMETRY ON VOID GROWTH IN DUCTILE MATERIALS WITH REALISTIC POROUS MICROSTRUCTURES

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<u>Summary</u> We have investigated the effect of plastic anisotropy and tension-compression asymmetry on void growth from materials which contain realistic porous microstructures considering 3D unit cells and two approaches: crystal plasticity models and macroscopic models. For that purpose, we have performed finite element calculations of cubic unit-cells which are subjected to periodic boundary conditions and include porosity distributions representative of different additively manufactured metals.

#### **INTRODUCTION**

In order to elucidate the role of real void sizes or real intervoid distances on the mechanisms of ductile fracture, finite element simulations, including distributions of voids obtained from X-ray tomography analysis of porous materials, bring about opportunities to study ductile damage in realistic situations. The main objective of our work is to analyse the effect of plastic anisotropy and tension-compression-asymmetry on void growth in ductile materials with realistic porous microstructures. The matrix material will be described considering two approaches: a crystal plasticity framework [1] and classic anisotropic formulations [2],[3] and [4]. To the authors' knowledge, an experimentally-based void configuration with a number of pores large enough to ensure a significant statistical representation of the porous microstructure was never fully mapped within a 3D representative volume element.

#### CONSTITUTIVE MODELS Stewart and Cazacu (2006) 's macroscopic model [5]:

$\mathbf{F} = \mathbf{F}^{\mathbf{e}} \mathbf{F}^{\mathbf{p}}.$	$\dot{\hat{ au}}_{lpha} = h_0 \sum_{eta=1}^{N^S} q_{lphaeta} \left(1 - \frac{\hat{ au}_{eta}}{\hat{ au}_s}\right)^{P_2}  \hat{\gamma}_{eta} $	$f = \tilde{a}$	$\dot{r} - \sigma_1$		ō	= ( \$\phi\$	CPB06	$\int_{0}^{1/a} \phi_{C}$	PB06 = 5	$ \tilde{s}'_i  -$	$k\tilde{s'_i})^a$
$\tilde{\mathbf{F}}^{p} = \left(\sum_{\alpha=1}^{N^{S}} \dot{\hat{\gamma}}_{\alpha} \widetilde{\mathbf{M}}_{\alpha}\right) \mathbf{F}^{p},$	$\tau_{\mathbf{x}} = \widetilde{\mathbf{T}} \cdot \widetilde{\mathbf{M}}_{\mathbf{x}}$		[La	Luz	Liz	0	η 0	01	ī	=1	
$\mathbf{n}(t) = \left(\mathbf{F}^{\mathbf{e}\mathbf{T}}\right)^{-1} \tilde{\mathbf{n}}_{\mathbf{v}}$	$\overset{\sim}{M}_{\alpha}=\overset{\sim}{m}_{\alpha}\otimes \overset{\sim}{n}_{\alpha} \overset{\sim}{T}=\left(F^{eT}F^{e}\right)\overset{\sim}{S}\approx F^{p}SF^{pT}$		$L_{12}$	$L_{22} \\ L_{23}$	L <sub>23</sub> L <sub>23</sub>	0	0	0			
$ \tau_{\alpha} ^{p_1}$		L =	0	0	0	$L_{44} = 0$	$\begin{array}{c} 0 \\ L_{55} \end{array}$	0			
$\gamma_{\alpha} = \gamma_0 \left  \hat{\tau}_{\alpha} \right   \operatorname{sign}(\tau_{\alpha}).$	$\mathbf{m}(t)_{\alpha} = \mathbf{F}^{\mathbf{c}} \mathbf{m}_{\alpha},  \mathbf{S} = \frac{\partial \mathbf{I}}{\partial \mathbf{E}^{\mathbf{c}}} = \mathbb{C} \cdot \mathbf{E}^{\mathbf{c}}$		0	0	0	0	0	L <sub>66</sub>			

FINITE ELEMENT MODEL



Figure 1. (a) semi-transparent view displaying the porous microstructure, (b) loading conditions with  $\Sigma_1$ ,  $\Sigma_2$  and  $\Sigma_3$  being the principal values of the macroscopic stress tensor and (c) cut-view showing the fine mesh around the voids [2].

#### SAMPLE RESULTS



Figure 2. Distribution of misorientation angle with respect to the initial orientation for softhard bi-crystal for T = 0, L = -1 and equivalent strain  $\overline{\epsilon} = 0.45$ . Pole figures 111 present the spread of the current crystal orientations for each Gauss point in the cell [1]

The finite element model (see Fig. 1) used in simulations is a cubic unit-cell containing spatially distributed spherical voids of different sizes. This cell model is considered to be a representative volume element of a porous material. The equations for the nodal displacements reported in [1] have been used to impose periodic boundary conditions on the unit-cell, so that the displacement of opposed external nodes is coupled. Constant and controlled values of the macroscopic stress triaxiality and Lode parameter are prescribed during the entire loading history.

**Crystal plasticity model:** 

Figure 2 presents the influence of void on the heterogeneity of lattice rotation in a soft-hard bi-crystal, at T = 0 and L=1. Initially all elements within each half cell have the same orientation. It becomes apparent that while for a bi-crystal without a void, lattice rotates uniformly within each half cell, strong variation is observed when the voids are present at the grain boundary. The largest misorientation angle (> 45) is observed close to the void boundary, where the void is distorted most.

Fig. 3 shows the evolution of f/f0 with the angular misalignment  $\theta$  for L=-1 and T=0.33, 1 and 2. Results correspond to aluminium alloy 2090-T3 for material orientations in the case of x $\Sigma$ III y $\theta$  z $\theta$ , which means the minor loading direction  $\Sigma$ III is parallel to the rolling direction x, and the transverse and normal orthotropy axes, y and z, form an angle  $\theta$  with the other two loading directions. The minimum is shallow for T = 0.33, and much stronger for T = 2, so that the nonlinearity of the f/f0- $\theta$ . curves become more important as the triaxiality increases, bringing out the interplay between angular misaligment and stress triaxiality on void growth. The angle  $\theta$  for which the minimum occurs, while hardly dependent on the value of T, depends strongly on the material orientation.



Fig.3. Influence of the macroscopic stress triaxiality on void growth for AA 2090-T3. Unit-cell finite element calculations for L = -1 and T = 0.33, 1 and 2. Evolution of the normalized void volume fraction f/f0 with the angular misalignment  $\theta$  and material orientation x $\Sigma$ III y $\theta$  z $\theta$  [3]

(a) (b)  $\frac{1}{2}$   $\frac{1}{2$ 

Fig. 4 presents contours of effective plastic strain for a porous microsture with 27 voids and 5 clusters. The cluster contains voids which grow and rapidly interact with each other shortly after the loading starts, such that most of the pores are no longer spherical for  $\varepsilon = 0.033$ , with the voids surrounding the largest pores displaying a mushroom shape with a flatten face corresponding to the formation of an intervoid ligament. The flattening of the voids generally starts earlier than in the case of the porous microstructures with randomly distributed voids since packing the voids into clusters decreases the distance between voids. For  $\varepsilon = 0.066$  the plastic strain in the surface of most pores is greater than 1, and the separation between voids is minimal, depicting the beginning of coalescence

Figure 4. Results corresponding to porous microstructure with 27 voids and 5 clusters. Stress triaxiality T = 3 and Lode parameter L = -1. Snapshots corresponding to different values of macroscopic effective strain: (a)  $\overline{\epsilon} = 0$ , (b)  $\overline{\epsilon} = 0.033$  and  $\overline{\epsilon} = 0.066$  [4].

#### SUMMARY AND CONCLUSIONS

We have carried out a comprehensive numerical investigation on the role of real porous microstructure on the ductile behavior or 3D printed metals. The matrix material was described considering two approaches: a crystal plasticity framework and classic anisotropic formulations. The simulations have been carried out with random spatial distributions of voids and with clusters of the same size but different void densities, and the results have been compared to unit-cells with a single central pore. Calculations exchanging the loading directions for a given distribution of void sizes and positions have been performed to elucidate whether the evolution of the porous microstructure leads to anisotropic behavior of the unit-cell.

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